

## ABSTRACT

Recent studies have revealed that CA is not necessarily only a Weyl-node property, but is rather a Fermi surface property, and is also present in a more general class of materials, for example, in spin-orbit-coupled noncentrosymmetric metals (SOC-NCMs). Using a tight-binding model for SOC-NCMs, we first demonstrate that strain in SOC-NCMs induces anisotropy in the spin-orbit coupling and generates an axial electric field. Then, using the quasi-classical Boltzmann transport formalism with momentum-dependent intraband and interband scattering processes, we show that strain in the presence of external magnetic field can generate temperature gradients via the Nernst-Ettingshausen effect, whose direction and behavior depends on the interplay of multiple factors: the angle between the applied strain and magnetic field, the presence of the chiral anomaly, the Lorentz force, and the strength of interband scattering.

## INTRODUCTION

- Chiral anomaly have been observed in SOC-NCMs, which renders it to be a Fermi surface property rather than a Weyl-node property.
- SOC-NCMs have an additional  $k^2$  dependence in the hamiltonian. As a result, there will be two fermi surfaces associated with a single node.
- Strain induces anisotropy in the spin-orbit coupling and generates an axial electric field.
- Strain in the presence of external magnetic field can generate temperature gradients via the Nernst-Ettingshausen effect whose direction and behavior depends on the angle between the applied strain and magnetic field, and the strength of interband scattering.

## MATHEMATICAL MODELLING

- For SOC-NCMs under strain, we consider the Hamiltonian:

$$\frac{\hbar^2 k^2}{2m} + \alpha \mathbf{k} \cdot \boldsymbol{\sigma} + b_z k_z \sigma_z$$

We have taken into account the effect of Berry curvature on the semi-classical dynamics of Bloch electrons as per the equations:

$$\dot{\mathbf{r}}^\lambda = D^\lambda \left( \frac{e}{\hbar} (\mathbf{E} \times \boldsymbol{\Omega}^\lambda + \frac{e}{\hbar} (\mathbf{v}^\lambda \cdot \boldsymbol{\Omega}^\lambda) \mathbf{B} + \mathbf{v}_k^\lambda) \right)$$

$$\dot{\mathbf{p}}^\lambda = -e D^\lambda (\mathbf{E} + \mathbf{v}_k^\lambda \times \mathbf{B} + \frac{e}{\hbar} (\mathbf{E} \cdot \mathbf{B}) \boldsymbol{\Omega}^\lambda)$$

- We solved the Maxwell-Boltzmann Equation for a given distribution function without loss of generality.

$$\frac{\partial f^\lambda}{\partial t} + \dot{\mathbf{r}}^\lambda \cdot \nabla_{\mathbf{r}} f^\lambda + \dot{\mathbf{k}}^\lambda \cdot \nabla_{\mathbf{k}} f^\lambda = I_{\text{coll}}[f^\lambda]$$

and in accordance with Fermi-Golden rule we have,

$$I_{\text{coll}}[f_{\mathbf{k}}^m] = \sum_p \sum_{\mathbf{k}'} W_{\mathbf{k},\mathbf{k}'}^{mp} (f_{\mathbf{k}'}^p - f_{\mathbf{k}}^m)$$

Where,

$$W_{\mathbf{k},\mathbf{k}'}^{mp} = \frac{2\pi}{\hbar} \frac{n}{V} |\langle \psi_{\mathbf{k}'}^p | U_{\mathbf{k}\mathbf{k}'}^{mp} | \psi_{\mathbf{k}}^m \rangle|^2 \delta(\epsilon_{\mathbf{k}'}^p - \epsilon_{\mathbf{k}}^m)$$

- In the presence of finite temperature, the Maxwell Boltzmann equation is modified as:

$$\left[ \left( \frac{\partial f_{\mathbf{k}}^\lambda}{\partial \epsilon_{\mathbf{k}}^\lambda} \right) \left( e\mathbf{E} + \frac{\epsilon_{\mathbf{k}}^\lambda - \mu}{T} \nabla T \right) \cdot \left( \mathbf{v}_{\mathbf{k}}^\lambda + \frac{e\mathbf{B}}{\hbar} (\boldsymbol{\Omega}^\lambda \cdot \mathbf{v}_{\mathbf{k}}^\lambda) \right) \right] = -\frac{1}{D^\lambda} \sum_{\lambda'} \sum_{\mathbf{k}'} W_{\mathbf{k}\mathbf{k}'}^{\lambda\lambda'} (g_{\mathbf{k}'}^{\lambda'} - g_{\mathbf{k}}^{\lambda})$$

- The distribution function is obtained by solving the above equation, and is substituted in the current density equation:

$$\mathbf{J} = -\frac{e}{V} \sum_{\lambda,k} \dot{\mathbf{r}}_{\mathbf{k}}^\lambda f_{\mathbf{k}}^\lambda$$

- From the above equation we get the Hall and Peltier coefficients driven by electric and temperature gradient responses.

- At the steady state, the current density vanishes and we get the following equation:

$$\frac{\nabla T}{E_5} = \hat{\alpha}^{-1} \hat{\sigma}$$

## RESULTS

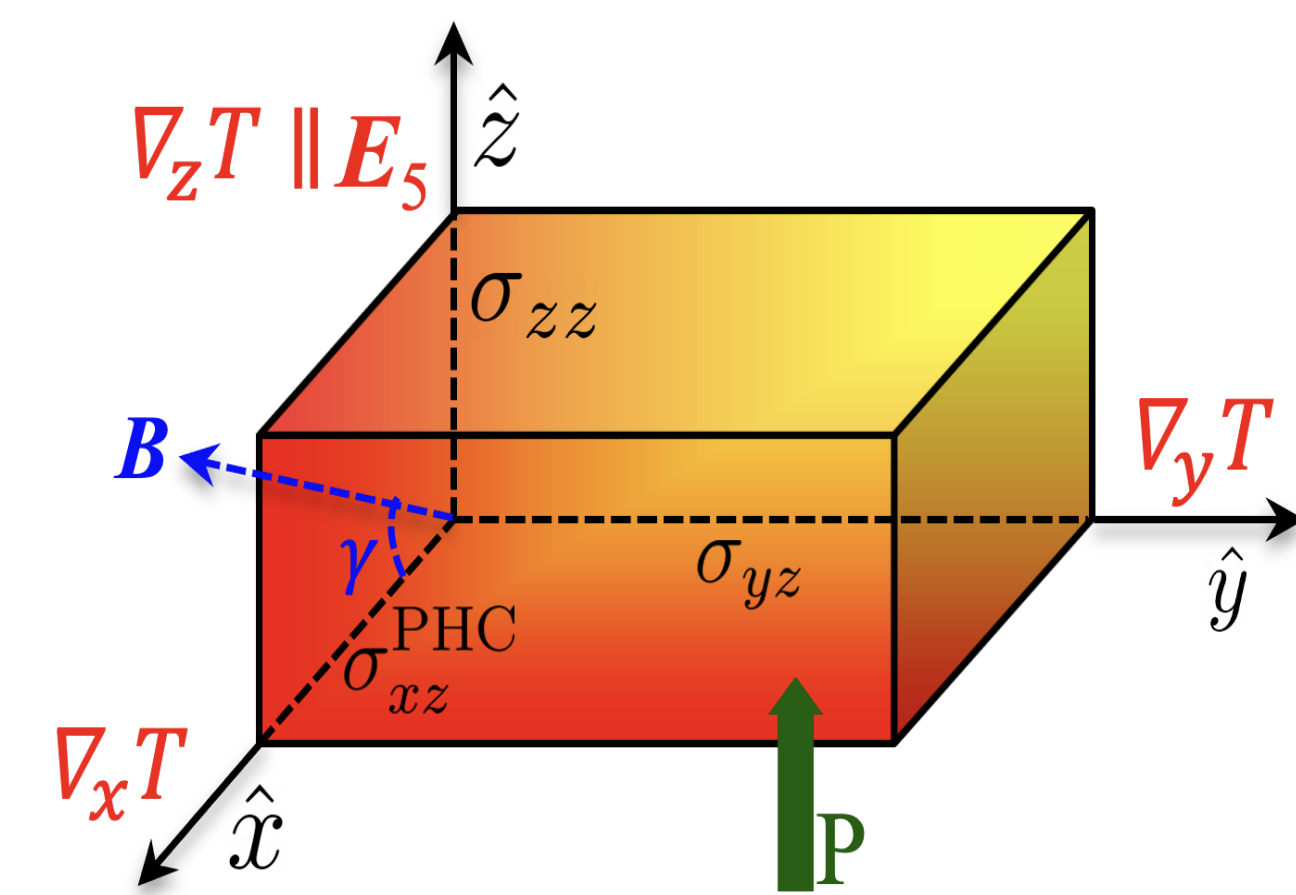


Fig1: Schematic that shows the strain-induced Ettingshausen effect

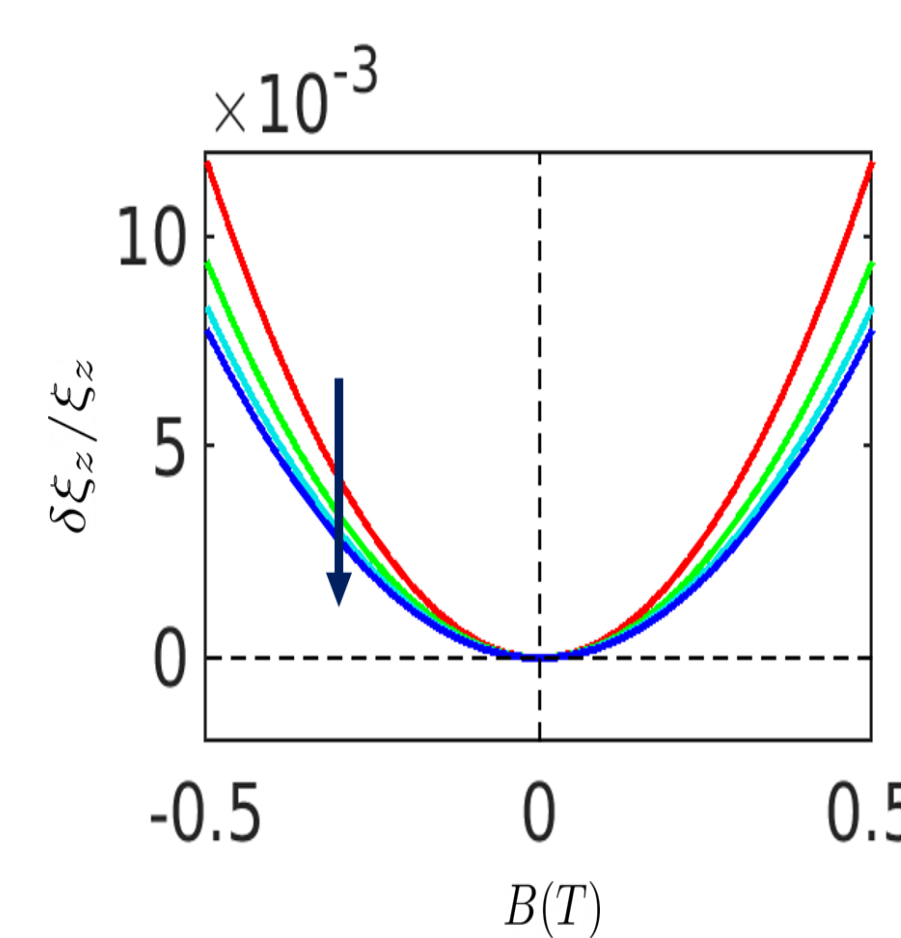


Fig 2: Temperature gradient for the case when direction of strain and magnetic field are parallel.

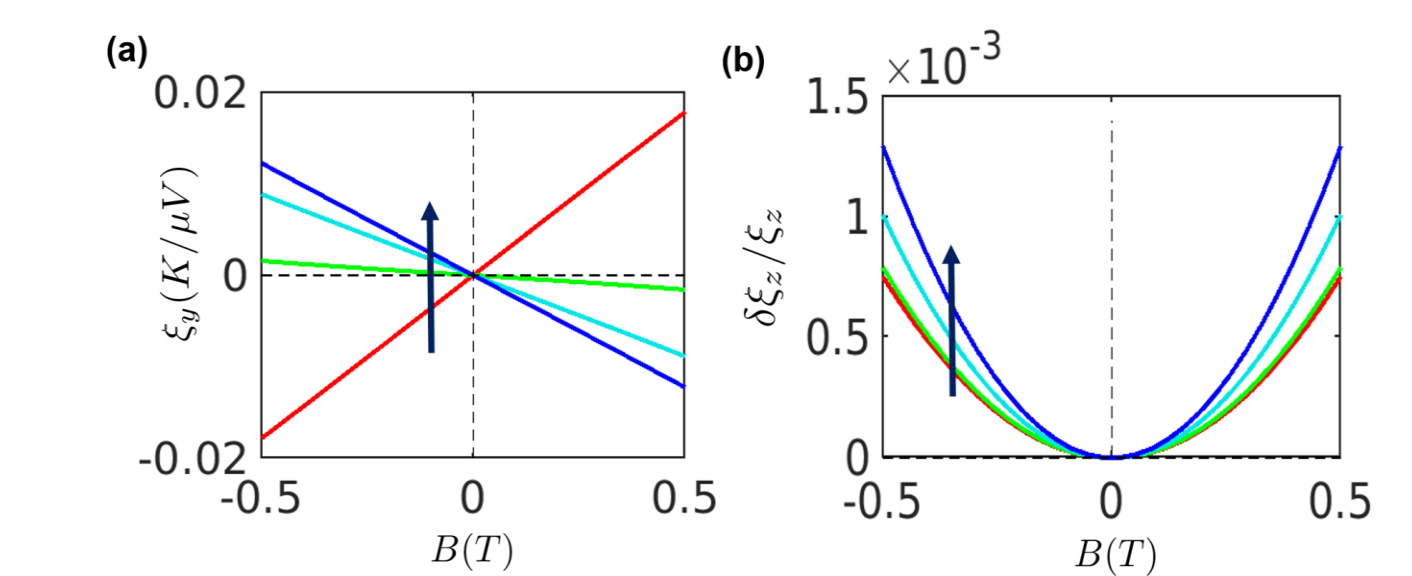


Fig 3: Temperature gradients for the case when strain and magnetic field are perpendicular.

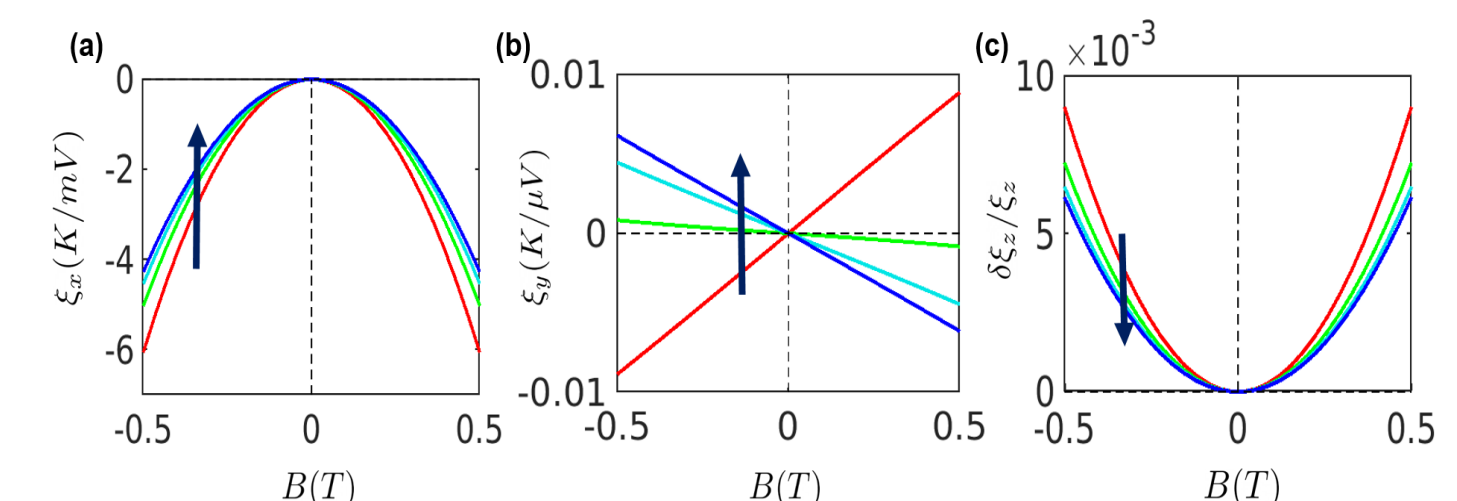


Fig 4: temperature gradient for an arbitrary angle.

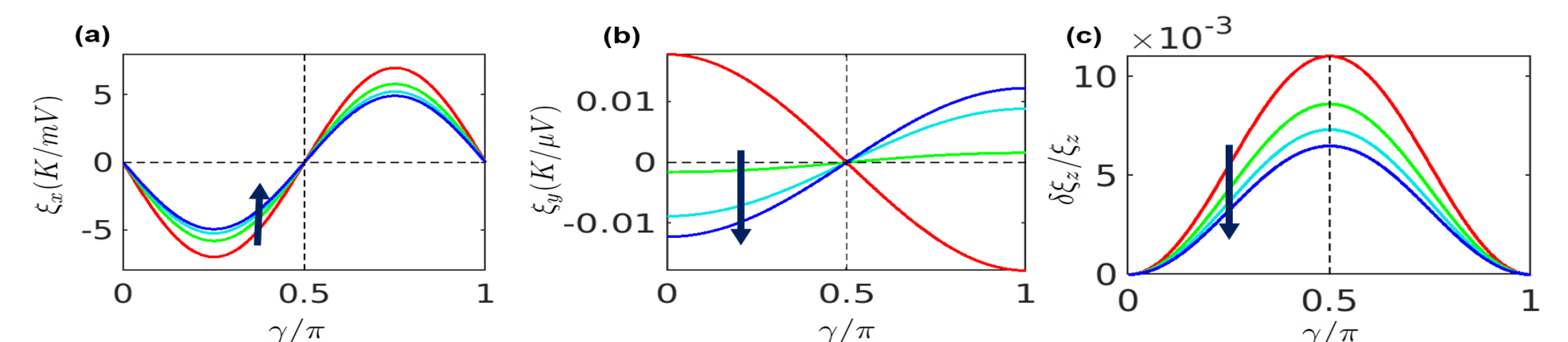


Fig 5: Dependence of temperature gradients on the angle between strain direction and magnetic field.

## CONCLUSIONS

- SOC-NCMs will have anisotropy in the spin-orbit coupling due to strain.
- Strain generates an axial electric field in SOC-NCMs along the direction of application.
- At the steady state, strain generates temperature gradients whose direction depends on the angle between the strain direction and magnetic field and the strength of interband scattering.
- Temperature gradient along x-direction is purely induced by chiral anomaly via PHC, while that along the y- and z-directions are generated by Lorentz Hall and LMC (chiral anomaly induced) respectively.

## REFERENCES

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